

BEYOND POPULAR SCIENCE



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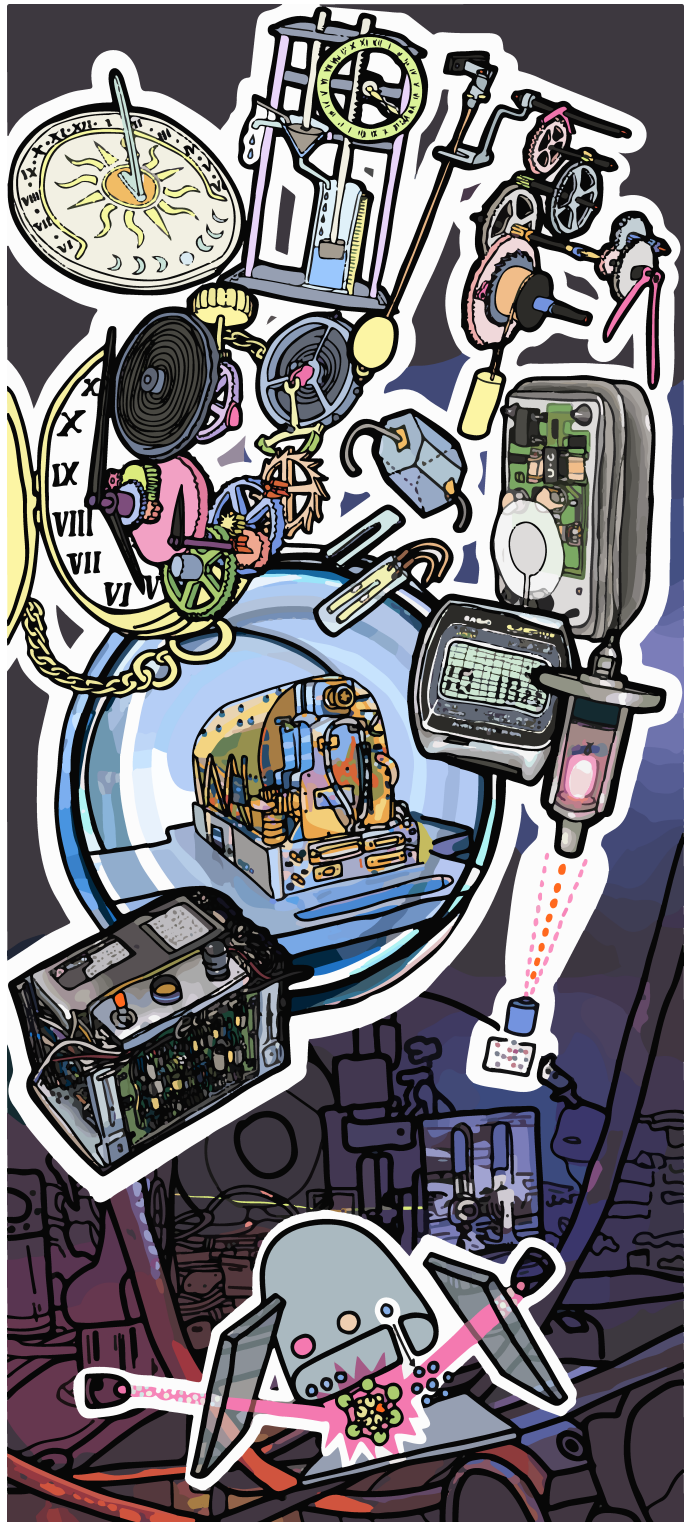
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**Timing Is
Everything**

Top (Mechanical and Early Timekeeping): The image begins with sundials—using cast shadows to track solar position. Mechanical clocks follow: weight-driven escapements, pendulums, and balance springs. These systems introduced regular, countable oscillations, enabling precision independent of sunlight.

Middle (Electronic and Atomic Clocks): The transition to electronic timing brought quartz crystal oscillators, tuning forks, and circuit-based digital clocks. Atomic clocks, such as the cesium beam and rubidium standards, use microwave transitions in atoms to define the second with extreme precision. These devices revolutionised navigation, telecommunications, and metrology.

Bottom (Next-Generation Nuclear Clocks): The current frontier involves optical lattice and nuclear clocks—probing energy levels in nuclei rather than electrons. These offer unprecedented temporal stability and resolution. If perfected, they could detect minute variations in gravity, test general relativity, and redefine the basic unit of time.



Timing Is Everything

Timekeeping has progressively moved toward smaller physical phenomena: from Earth's rotation to pendulums, from crystal oscillations to atomic transitions, and now toward nuclear resonances. The SI second, defined by 9,192,631,770 periods of cesium-133's hyperfine transition, relies on quantum interactions between nuclear and electronic magnetic moments. This shift to microscopic reference standards improves precision exponentially—hydrogen masers achieve stability of one part in 10^{13} , while optical lattice clocks using strontium reach one part in 10^{18} by probing transitions at $\sim 10^{15}$ Hz. The progression continues toward nuclear clocks using thorium-229, which promises precision of one part in 10^{19} by exploiting transitions in atomic nuclei rather than electron shells.



TIME AS COORDINATE ◦ CAESIUM-133
STANDARD ◦ 9192631770 Hz ◦ HYPERFINE
TRANSITIONS ◦ OPTICAL LATTICE CLOCKS ◦ FREQUENCY
COMBS ◦ GPS TIME CORRECTIONS ◦ THORIUM-229 NUCLEAR
CLOCK ◦ CHRONOMETRIC GEODESY ◦ PROPER TIME
INVARIANCE ◦ FUNDAMENTAL PHYSICS PROBE

„Zeit ist das, was man an der Uhr abliest.“

(“Time is what a clock measures.”)

— Albert Einstein, 1926

Timing Is Everything

Era / Technology	Accuracy	Size (m)	Time Reference
Sundials (1800 BCE)	10^{-2}	10^1	Solar shadow on gnomon
Ancient Water Clocks	10^{-3}	10^{-1}	Liquid level change
Verge Clocks (13th c.)	10^{-2}	10^{-1}	Crown wheel / verge foliot
Pendulum Clocks (1656)	10^{-5}	10^0	Pendulum arc length
Marine Chronometer (18th c.)	10^{-6}	10^{-2}	Balance spring oscillator
Quartz Oscillators (1930s)	10^{-8}	10^{-3}	Crystal thickness (MHz mode)
Ammonia Maser (1953)	10^{-9}	10^{-10}	NH_3 inversion barrier
Cesium Beam Standard (1955)	10^{-10}	10^{-10}	^{133}Cs hyperfine structure
Hydrogen Maser (1960s)	10^{-13}	10^{-10}	^1H hyperfine structure
Rubidium Vapour	10^{-11}	10^{-10}	^{87}Rb hyperfine structure
Cesium Fountain (1990s)	10^{-15}	10^{-10}	Interference of free atoms
Optical Lattice (2010s)	10^{-18}	10^{-10}	Atomic dipole transitions
Projected Thorium Nuclear	10^{-20}	10^{-14}	Intrinsic nuclear excitation

Time is a coordinate assigned to events, an ordering imposed on phenomena, and a physical quantity whose measurement depends on the reproducibility of periodic processes. The challenge in defining time stems from its dual character: operationally, time is what clocks measure, but physically, clocks are systems that embody time through the regularity of their transitions. A theory of time must therefore address both its measurement and its assignment.

Historically, time was defined by external reference. A day was one full rotation of the Earth, a year one revolution around the Sun. These intervals were directly observable but not uniform. Earth's rotation slows due to tidal friction, and its orbit varies minutely from year to year. As clocks improved, it became clear that astronomical cycles were not perfectly periodic. They were also not universally accessible.

Modern definitions turn inward. Time is now anchored to the internal configuration of matter. A clock is a system that undergoes periodic change—a pendulum, a quartz crystal, or a quantum oscillator—and time is defined by counting these cycles. The second is defined by the *Système international d'unités* (SI). The SI second is defined as exactly 9,192,631,770 periods of the hyperfine transition of the cesium-133 atom.

Still, the concept of time requires consistency. If time is relative—as in special and general relativity—how can clocks agree? The answer lies in local invariance and synchronisation protocols. In special relativity, time intervals are frame-dependent, but proper time—the time measured along an observer's worldline—is invariant. In general relativity, the curvature of spacetime causes time to flow at different rates in different gravitational potentials. Atomic clocks confirm that indeed identical devices tick faster at altitude than at sea level. Yet these variations are predictable and correctable.

To coordinate time across systems and locations, one defines a reference frame and applies relativistic corrections. Global time standards, such as International Atomic Time (TAI), are constructed by ensemble averaging signals from many atomic clocks, each corrected

for gravitational potential and velocity. Clocks that are precise but prone to long-term drift are periodically resynced against clocks that are more accurate; the ensemble combines the strengths of both. The result is a global timescale without assertion of universal time.

Time also enters theoretical physics as a parameter. In Newtonian mechanics, time is absolute and flows uniformly. In quantum mechanics, it appears as an external parameter in the Schrödinger equation (Schrödinger, 1926). In general relativity, time is a coordinate entangled with space, whose flow is determined by the metric tensor. Time serves as an index that parameterizes change.

In quantum field theory and statistical mechanics, time appears asymmetrically. The microscopic laws are time-reversal symmetric, yet macroscopic systems exhibit irreversibility. The asymmetry is imposed not by the equations, but by boundary conditions and coarse-graining (the replacement of a detailed description with a statistical one). The direction of time—the arrow from past to future—emerges from the configuration of initial conditions and the growth of entropy (Boltzmann, 1877).

The measurement principle is that time is a relation between events measured by clocks as intervals. The definition principle is that the flow of a system is the unfolding of configurations in accordance with dynamical laws that govern evolution from past to future.

Physical timekeeping builds upon these principles. The invariance of atomic transitions allows time to be physically instantiated as a countable quantity, realised through interactions with matter that exhibit extraordinary regularity. Atomic clocks operationalize time by coupling electromagnetic fields to well-defined quantum transitions—processes governed by the internal energy levels of atoms. These transitions occur at precise frequencies determined by the laws of quantum electrodynamics and the values of fundamental constants, making them immune to most environmental and instrumental variations. The resulting periodicity is intrinsic.

In the case of cesium-133, the phenomenon that defines the second is the hyperfine splitting of its ground electronic state. The splitting arises from the interaction between two magnetic moments: that of the nucleus, which acts as a tiny bar magnet due to its intrinsic spin, and that of the valence electron, whose magnetic field is generated by both its orbital motion and its intrinsic spin. These moments couple through the magnetic dipole interaction, producing a small energy difference between two configurations. Quantum mechanically, the total angular momentum of the atom is given by $\vec{F} = \vec{I} + \vec{J}$, where \vec{I} is the nuclear spin and \vec{J} the total electronic angular momentum. In cesium-133, which has nuclear spin $I = 7/2$ and electronic angular momentum $J = 1/2$ in its ground state, this coupling results in two hyperfine levels: $F = 4$ and $F = 3$.

The transition between these levels occurs at a microwave frequency of 9.192631770 GHz. Because this energy difference is sharply defined and identical for all cesium-133 atoms in isolation, it serves as a natural frequency reference. The transition is measured by subjecting a cloud of cesium atoms to a tunable microwave field while monitoring population redistribution between the two states. When the applied frequency matches the energy gap—satisfying the resonance condition $E = h\nu$ —atoms undergo induced

transitions, which can be detected via state-selective fluorescence or ionisation. In practice, a feedback loop adjusts the microwave oscillator to maximise this transition probability. The resulting frequency is then divided electronically to produce the one-second interval. The process defines the second as the number of cycles of this specific atomic transition.

Two distinct qualities characterise clock performance: *precision* (or stability) measures how consistently a clock reproduces the same interval from one tick to the next, while *accuracy* measures how closely that interval matches the true definition of the second. A clock can be extraordinarily precise—ticking with negligible variation over hours—yet slowly drift from the correct frequency over months. Conversely, a clock can be accurate on average but noisy tick-to-tick. The best timekeeping systems must excel at both.

Hydrogen masers generate coherent microwave radiation at 1.42 GHz via stimulated emission between hyperfine levels of atomic hydrogen. Their short-term frequency stability, driven by long coherence times in a wall-coated storage bulb, surpasses that of most other clock types. Although long-term drift limits their use as absolute standards, they serve as exceptional flywheel oscillators in timekeeping ensembles, bridging intervals between recalibrations from more accurate devices.

Rubidium clocks—especially chip-scale atomic clocks (CSACs)—offer compact, energy-efficient timing solutions for portable and embedded applications. These systems exploit optical pumping to polarise a vapour of ^{87}Rb atoms and monitor resonant microwave transitions via changes in transmitted light. The clock output disciplines an internal quartz oscillator, yielding fractional stabilities on the order of 10^{-11} to 10^{-12} , sufficient for GPS receivers, telecommunications, and low-power navigation.

Optical lattice clocks improve precision by probing narrow-linewidth electronic transitions in neutral atoms confined within standing-wave laser fields. At the ‘magic wavelength’ (species-dependent), the differential AC Stark shift between clock states vanishes (Katori, 1999), preserving the transition frequency despite optical confinement. Atoms such as strontium and ytterbium offer transition frequencies near 10^{15} Hz, and interrogation times exceeding one second yield quality factors above 10^{17} . These systems achieve fractional instabilities below 10^{-18} . Optical frequency combs enable comparison to microwave references, bridging domains and facilitating global synchronisation.

With such precision, relativistic effects become measurable and essential. Identical clocks placed at different gravitational potentials accumulate proper time at different rates due to gravitational redshift (Pound & Rebka, 1960). The shift $\Delta f/f = gh/c^2$ enables vertical positioning to centimetre resolution—the basis of chronometric geodesy. GPS satellites, which orbit at 20,200 km, exhibit both special relativistic time dilation (from orbital velocity) and gravitational blueshift (from altitude). Pre-launch frequency offsets and onboard corrections account for the net gain of approximately 38 microseconds per day, maintaining sub-metre positional accuracy.

Nuclear clocks aim to surpass atomic standards by exploiting transitions in the atomic nucleus, which are orders of magnitude less sensitive to electric and magnetic perturbations. The thorium-229 isomer exhibits the lowest known nuclear excitation energy—approximately 8.3 eV—placing it within reach of laser spectroscopy in the vacuum ultraviolet. Its

long radiative lifetime implies a millihertz-scale natural linewidth, suggesting a potential quality factor above 10^{19} . Two architectures dominate experimental development. In ion-trap systems, individual $^{229}\text{Th}^{3+}$ ions are confined by radiofrequency fields, laser-cooled, and interrogated using high-resolution VUV (vacuum ultraviolet) frequency combs. In the solid-state approach, thorium nuclei are embedded in wide-bandgap optical crystals such as CaF_2 or MgF_2 . These hosts suppress internal conversion decay and enable parallel interrogation of large ensembles. Challenges include spectral broadening from lattice inhomogeneity, background fluorescence, and the engineering of narrowband, stable VUV sources. Recent experiments have reported increasingly precise energy determinations and quantum-resolved spectroscopy of the transition, with rapid progress toward routine laser control.

The implications of nuclear timekeeping extend beyond metrology. Due to the fine balance of nuclear forces, the ^{229}Th isomer is predicted to be hypersensitive to variations in the fine-structure constant, scalar field couplings, or violations of local position invariance. Networks of synchronised thorium clocks could detect transient dark matter interactions or topological defects via correlated frequency excursions—simultaneous shifts in transition frequency caused by passing field disturbances that temporarily alter the local values of fundamental constants. Timekeeping becomes a probe of fundamental physics.

What was once derived from the rotation of celestial bodies is now defined by invariant atomic structure—and may soon be defined by the nucleus, whose internal dynamics offer a new frontier for precision and for discovery.

The Seven SI Base Units

Quantity	Unit (Symbol)	Definition
Time	second (s)	9,192,631,770 periods of the cesium-133 hyperfine transition
Length	metre (m)	Distance light travels in vacuum in $1/299,792,458$ of a second
Mass	kilogram (kg)	Fixed by setting Planck's constant $h = 6.62607015 \times 10^{-34}$ J·s
E. Current	ampere (A)	Fixed by setting elementary charge $e = 1.602176634 \times 10^{-19}$ C
Temperature	kelvin (K)	Fixed by setting Boltzmann constant $k_B = 1.380649 \times 10^{-23}$ J/K
Amount	mole (mol)	Exactly $6.02214076 \times 10^{23}$ elementary entities
Lum. Intensity	candela (cd)	Fixed by setting $K_{cd} = 683$ lm/W at 540 THz

The 2019 redefinition marked the final transition from artefact-based standards—such as the kilogram prototype stored in Paris—to definitions anchored in immutable fundamental constants, ensuring universal reproducibility without dependence on physical objects that degrade, drift, or require secure storage.

Exercises

Fusible Numbers: Exercises in Constructive Time

A classic riddle: given two candles, which each burn for 1 h, how can you measure 45 m?

Fusible numbers form a well-ordered subset of the rationals constructed iteratively from zero. A number z is fusible if there exist previously constructed fusible numbers x and y , with $|x - y| < 1$, such that $z = (x + y + 1)/2$. The construction corresponds to lighting a unit-time fuse at both ends with delay. Results about the growth of certain associated functions are linked to very fast-growing hierarchies and have connections to independence from Peano arithmetic.

1. The Fuse Construction

For a unit fuse lit at time x on one end and time y on the other (with $|x - y| < 1$), prove it extinguishes at $z = (x + y + 1)/2$. Consider the burn dynamics when both flames are active.

2. Enumeration Below 2

Determine all fusible numbers < 2 by systematic application of the construction rule.

3. Dyadic Structure

Prove that all fusible numbers have the form $a/2^k$ for integers $a \geq 0$ and $k \geq 0$.

4. The Margin Function

Let a_n be the smallest fusible number exceeding n . Prove that

$$a_n = n + \frac{1}{2^{k(n)}}$$

for some $k(n) \in \mathbb{N}$. Compute $k(0)$, $k(1)$, $k(2)$. The function $k(n)$ grows extremely rapidly; certain asymptotic properties are connected to statements independent of Peano arithmetic.

5. Well-Ordering

Prove that the fusible numbers form a well-ordered subset of \mathbb{Q}^+ . What does this imply about infinite decreasing sequences?

Context

Fusible numbers (Erickson, Xu) demonstrate how elementary constructions yield extremely fast growth. The margin values are:

$$a_0 = \frac{1}{2}, \quad a_1 = 1 + \frac{1}{8}, \quad a_2 = 2 + \frac{1}{1024}.$$

The growth behaviour connects to ordinal ε_0 . (See Chapter 17)

Quantum Transitions and the Limits of Clock Stability

Atomic and nuclear clocks define time by referencing a sharply resonant transition between two quantum states. The frequency of this transition is determined by fundamental constants and is reproducible across identical systems. The precision with which this frequency can be measured depends on the linewidth of the transition, the stability of the interrogation system, and the protocol used to extract frequency information. This section formalises the mathematical quantities that govern frequency stability, relates them to the physical structure of the transition, and identifies the limits imposed by spacetime curvature and coupling constants.

Spectral Linewidth and Quality Factor

Let f_0 denote the central transition frequency and Δf the full width at half maximum (FWHM). The quality factor is defined by:

$$Q = \frac{f_0}{\Delta f}.$$

In Ramsey interrogation, $\Delta f \approx 1/(2T)$, where T is the free evolution time between pulses. Hence,

$$Q \approx 2f_0T.$$

Optical lattice clocks probing transitions in strontium or ytterbium atoms with $f_0 \sim 10^{15}$ Hz and $T \sim 1$ s routinely achieve $Q > 10^{15}$.

Allan Deviation and Averaging Behaviour

The fractional instability of a clock over averaging time τ is quantified by the Allan deviation:

$$\sigma_y(\tau) \approx \frac{1}{Q} \cdot \frac{1}{\text{SNR}} \cdot \sqrt{\frac{T_c}{\tau}},$$

where SNR is the signal-to-noise ratio and T_c the cycle time. Increasing Q , improving detection fidelity, and lengthening τ all contribute to reduced $\sigma_y(\tau)$.

Hyperfine and Nuclear Transition Energies

In cesium-133, the clock transition arises from magnetic dipole coupling between nuclear spin \vec{I} and electron angular momentum \vec{J} , producing total angular momentum $\vec{F} = \vec{I} + \vec{J}$ and energy splitting:

$$E_F = \frac{A}{2} [F(F+1) - I(I+1) - J(J+1)].$$

The $F = 3 \leftrightarrow F = 4$ transition at $f_0 = 9.192\,631\,770$ GHz defines the SI second. In ^{229}Th , the nuclear excitation energy $E \approx 8.3$ eV corresponds to:

$$f_{\text{Th}} = \frac{E}{h} \approx 2.0 \times 10^{15} \text{ Hz}, \quad Q_{\text{Th}} \gtrsim 10^{19}.$$

Relativistic Shift and Coupling Sensitivity

In general relativity, clocks at different gravitational potentials W accumulate proper time at different rates. The fractional frequency shift is:

$$\frac{\Delta f}{f} = \frac{\Delta W}{c^2} \approx \frac{gh}{c^2},$$

where g is gravitational acceleration and h the height difference. At 10^{-18} resolution, height differences of 1 cm are resolvable.

Clock transitions sensitive to the fine-structure constant α respond to coupling variations via:

$$\frac{\Delta f}{f} = K_\alpha \cdot \frac{\Delta \alpha}{\alpha},$$

where K_α is a dimensionless sensitivity coefficient. In nuclear systems such as ^{229}Th , this coefficient may exceed 10^4 , amplifying the clock's utility in probing scalar fields or dark sector interactions.

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